## Letters to the Editor

## **Concentration-Dependence of Nonelectrolyte Permeability of Toad Bladder**

I should like to offer the following comments on Chen and Walser's recent paper [2]. In Appendix A, presented as a mathematical proof of Eq. (12), they derive a relationship between the unidirectional flux of a nonelectrolyte in the absence of net flow, measured under two circumstances ( $J_{\pm}$ , with the mucosal-serosal concentration difference  $\Delta c \equiv c_m - c_s$  and hydrostatic pressure difference  $\Delta p$  both equal to zero, and  $J_{\pm}^{eq}$ , with  $\Delta p = \Delta p_{eq}$ , the value appropriate to compensate for a given value of  $\Delta c \pm 0$ ). Reasoning in terms of a Taylor's expansion, they write

$$J_{\leftrightarrows}^{\mathrm{eq}}/J_{\leftrightarrows} = 1 + \sum_{k=1}^{j} a_k (c_m - c_s)^k.$$
(A1)

This formulation ignores the consideration that in principle the unidirectional flux is a function of hydrostatic pressure as well as concentrations, so that a complete Taylor series comprises terms in both  $\Delta c$  and  $\Delta p$ , as well as mixed terms.

Leaving this point aside, for the case of constant  $c_s$ , Eq. (1) is transformed into

$$J_{\ddagger}^{eq}/J_{\ddagger} = 1 + \xi_1 [(c_m - c_s)/c_s] + \xi_1 [(c_m - c_s)/c_s]^2 + \xi_1 [(c_m - c_s)/c_s]^3 + \dots$$
(A3)

where  $\xi_1 = a_k c_s^k$ . Stating that the  $\xi_1 s$  are binomial coefficients, it is then concluded that

$$J_{\ddagger}^{\text{eq}}/J_{\ddagger} = [1 + (c_m - c_s)/c_s]^{\xi_1} = (c_m/c_s)^{\xi_1}$$
(A4); (A5)

where  $\xi_1 = a_1 c_s$ . (The denominator  $J_{\pm}$  in Eq. (A4) was omitted.)

In analyzing the above, it is unclear why it is considered that the  $\xi_1 s$  are binominal coefficients since, even on the basis of Eq. (A1) as written,  $\xi_1$  is a complicated quantity, being a function of k-th order derivatives of the unidirectional fluxes with respect to concentration, which in principle might be expected to depend on concentrations, hydrostatic pressures, and membrane parameters. On the other hand, a binomial coefficient  $C_k^{\xi_1} = \xi_1(\xi_1 - 1)$   $(\xi_1 - 2)...(\xi_1 - k + 1)/k!$ , and there is no reason to expect in general that  $\xi_1$  need equal  $C_k^{\xi_1}$ .

For these reasons, I do not feel that Appendix A establishes the validity of Eq. (A5) or the relationships deduced from it. In particular, it appears that, as previously, Eq. (12) must be regarded as a postulate [1].

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#### References

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# **Reply to: Concentration-Dependence** of Nonelectrolyte Permeability of Toad Bladder

In reply to Dr. Essig, we agree that in the case of nonelectrolyte transport at constant temperature, the unidirectional solute flux is a function of gradients of hydrostatic pressure  $(\Delta p)$  and solute concentration  $(\Delta C)$ . However, it must be kept in mind that we are dealing only with passive transport systems at equilibrium. Under this condition, J=0 in which J represents the net solute flux transported, and as described by Eq.(7) of our paper [1],

$$\Delta p = (RT/\overline{\nu}) \ln(C_{\rm m}/C_{\rm s}). \tag{7}$$

Hence, even if  $J_{=}^{eq}$  is a function of both  $\Delta p$  and  $\Delta C$ , since  $\Delta p$  is a function of  $C_m$ , according to Eq. (7), it is clear that  $J_{=}^{eq}$  is a function of the one independent variable,  $C_m$  [2]. To see this, we consider the unidirectional solute flux as a function of  $\Delta p$  and  $\Delta C$ . Note that this statement is also valid when the transport system is at equilibrium. Thus,

$$J_{\pm}^{\text{eq}} = J_{\pm}^{\text{eq}}(\Delta p, \Delta C). \tag{B1}$$

By the use of chain rules [2], we obtain from Eq. (B1)

$$dJ_{\pm\pm}^{\rm eq} = (\partial J_{\pm\pm}^{\rm eq} / \partial \Delta p)_{AC} d\Delta p + (\partial J_{\pm\pm}^{\rm eq} / \partial \Delta C)_{Ap} d\Delta C. \tag{B2}$$

From Eq. (7), for fixed  $C_s$  and constant T,

$$d\Delta p = (RT/\bar{\nu}C_m)d\Delta C. \tag{B3}$$

Introducing Eq. (B3) into Eq. (B2), we obtain, by rearranging,

$$dJ_{\Xi}^{eq} = \{ (RT/\bar{\nu} C_m) (\partial J_{\Xi}^{eq} / \partial \Delta p)_{\Delta C} + (\partial J_{\Xi}^{eq} / \partial \Delta C)_{\Delta p} \} d\Delta C$$
  
=  $f (\Delta C) d\Delta C$ , (B4)

which by integration over the thickness of the membrane gives the expression for  $J_{\pm 5}^{eq}$  in terms of  $C_m - C_s$ .

Dr. Essig also questions that in Eq. (A3)

$$J_{\frac{sq}{ss}}^{eq}/J_{\frac{sq}{ss}} = 1 + \zeta_1 \{ (C_m - C_s)/C_s \} + \zeta_1 \{ (C_m - C_s)/C_s \}^2 + \dots + \zeta_1 \{ (C_m - C_s)/C_s \}^k + \dots$$
(A3)

there is no reason to claim that the coefficient  $\zeta_1$  as

defined by  $\zeta_1 = a_k C_s^k$  need equal the binomial coefficient,  $C_k \zeta_1$ , because the parameter  $\zeta_1$  might be expected to depend on  $\Delta C$ ,  $\Delta p$  and membrane parameters. It should be remarked here that in writing Eq. (A1), which has been justified above to be a valid statement for transport systems at equilibrium, the coefficient  $a_k$ , in the Taylor's series must be evaluated by differentiating both sides of Eq. (A1) k times and setting  $C_m = C_s$ , i.e.,

$$a_k = J^{eq(k)}_{\varsigma}(C_s)/k! J_{\varsigma} \tag{B5}$$

where  $J_{i=1}^{eq(k)}(C_s)$  is the  $k^{th}$  derivative of  $J_{i=1}^{eq}$  evaluated at  $C_m = C_s$ . From Eq. (B5), clearly,  $a_k$  and thus  $\zeta_1$  are independent of  $\Delta p$  and  $C_m$  but depend on the constant parameter  $C_s$  and the membrane. Moreover, according to Eqs. (A1) and (B5), if all derivatives of  $J_{i=1}^{eq}$  exist at  $C_m = C_s$ , it is apparent that  $J_{i=1}^{eq}$  can be expressed by a binomial series as represented by Eq.(A3), since the coefficients  $(a_1, a_2...a_k)$  are not necessarily independent of each other [2].

Based on the above theoretical analysis, we conclude that the method used for the derivation of Eq. (12) as shown in Appendix A of our paper [1] is mathematically and physically justifiable and the validity of Eq. (A5) is thus warranted.

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